## Quantum Quasi-Monte Carlo

## for non-equilibrium quantum systems

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## Outline

I. Real time "diagrammatic" Quantum Monte Carlo. C. Bertrand, S. Florens, OP, X. Waintal Phys. Rev. X 9, 04I008 (2019)
Solution of the out of equilibrium quantum dot.

$$
\begin{aligned}
& \text { Phys. Rev. B 91, } 245154(2015) \\
& \text { Phys. Rev. B 100, } 125129(2019)
\end{aligned}
$$

## 2. Quantum Quasi-Monte Carlo.

How to compute the perturbative expansion faster ( $\mathrm{x} 100, \mathrm{xI} 0000$ ) and more precisely.

## Collaborators



Corentin Bertrand (Flatiron/CCQ)
Cf poster


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Bill Triggs (Grenoble, France)


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## Out of equilibrium \& strong correlations

- Many new experiments : Pump probe, quantum dots, ultra-cold atoms, cavities.


Pump probe


Nano-electronics


Ultra-cold atoms

- Computational physics challenge :
- Exact methods for out of equilibrium systems, at strong coupling
- Control, speed and precision
- Long time (after quench), steady state. Resolve various energy/time scales.


## Road map

Quantum dots


- Kondo effect.

L .Glazman et al. I988, P. Lee 1988.
D. Goldhaber-Gordon, 1998

- High precision benchmark in equilibrium (Bethe Ansatz)

Out of equilibrium
DMFT solvers


Lattice models
Real solid


- I atom+ self-consistent bath
- Few exact solvers
(inchworm Cf Cohen's talk, our work)
- Equilibrium : cf talk by M. Ferrero
- Non equilibrium.Transport.


## Perturbation theory

$$
Q(t)=\sum_{n=0}^{K} Q_{n}(t) U^{n}
$$

- Use perturbation theory $(\mathrm{K}=10-15)$, even deep in strong coupling regime (e.g. Kondo effect).
- Real time "diagrammatic" Quantum Monte Carlo (Cftalk of N. Prokof'ev, M. Ferrero)

How to compute $Q_{n}(t)$ ?
How to sum the series?

## How to compute $\mathrm{Q}_{\mathrm{n}}(\mathrm{t})$ ?

- Schwinger-Keldysh formalism $\mathrm{Q}_{\mathrm{n}}$ is a n -dimensional integral


$$
\begin{array}{r}
Q_{n}(t)=\frac{1}{n!} \int_{t_{0}}^{\infty} d u_{1} \ldots d u_{n}\left(\sum_{\alpha_{i}= \pm 1} \prod_{i} \alpha_{i} \operatorname{det}(\ldots)\right) \\
\equiv f_{n}\left(t, u_{1}, \ldots, u_{n}\right)
\end{array}
$$

Vertices.Times $u_{i}$.
Keldysh indices $\alpha=-I, I$

Profumo, Messio, OP, Waintal
PRB 91, 245I54 (2015)

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Profumo, Messio, OP, Waintal PRB 91, 245I54 (2015)
(Quasi) Monte Carlo Explicit sum

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Profumo, Messio, OP, Waintal PRB 91, 245154 (2015) (Quasi) Monte Carlo Explicit sum

- Long time limit $\mathrm{t} \rightarrow \infty$ is easy. $f_{n}$ is centered around t . Massive cancellations in the sum.
- No "dynamical sign problem" contrary to previous real time QMC, e.g. P.Werner et al PRB 2009
- $O\left(n^{3} 2^{n}\right)$ cost to compute $f_{n}(u)$. In practice, $n=10-15$.


## How to sum the series?



A finite radius of convergence! Singularities poles, branch cuts


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Profumo et al. PRB 91, 245154 (2015)
Bertrand et al. Phys. Rev. X 9, 041008 (2019)


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A finite radius of convergence ! Singularities poles, branch cuts

- Change of variable $\mathrm{W}(\mathrm{U})$, with $\mathrm{W}(0)=0$

$$
Q=\sum_{n \geq 0} Q_{n} U^{n}=\sum_{p \geq 0} \bar{Q}_{p} W^{p}
$$

A few results on the quantum dot

A simple model for the quantum dot

- Anderson model with two leads (L, R).


$$
\begin{gathered}
H=\sum_{\substack{k \sigma \\
\alpha=L, R}} \varepsilon_{k \alpha} c_{k \alpha}^{\dagger} c_{k \sigma \alpha}+\sum_{\sigma} \varepsilon_{d} d_{\sigma}^{\dagger} d_{\sigma}+U n_{d \uparrow} n_{d \downarrow}+\sum_{\substack{k \sigma \\
\alpha=L, R}} g_{k \sigma \alpha}\left(c_{k \sigma \alpha}^{\dagger} d_{\sigma}+\text { h.c. }\right) \\
\text { Bath } \quad \text { Dot Hybridization }
\end{gathered}
$$

- We want : current $\mathrm{I}\left(\mathrm{V}_{\mathrm{b}}\right)$, spectral function on the dot, Kondo effect, ...

Kondo effect in equilibrium

- Benchmark with NRG (numerical renormalisation group)

Spectral function on the dot


Kondo temperature


## Fermi liquid at low energy

- Equilibrium.

Self-energy, away from particle-hole symmetry

Bertrand et al. 2019
Phys. Rev. X 9, 041008 (2019)

Self energy (Re)

Self energy (Im)


## Out of equilibrium

- Destruction of the Kondo resonance by voltage bias


$$
T=0
$$



$$
T=\Gamma / 50
$$

- Particle hole asymmetric case



## Out of equilibrium distribution function of the dot

- Not a Fermi function

Bertrand et al. 2019
Phys. Rev. X 9, 041008 (2019)

- At $\mathrm{U}=0$ double step, due to the 2 Fermi leads.



# Quantum Quasi-Monte Carlo 

$\mathrm{Q}_{\mathrm{n}}(\mathrm{t})$ : a n -dimensional integral

$$
Q_{n}(t)=\frac{1}{n!} \int_{t_{0}}^{\infty} d u_{1} \ldots d u_{n} f_{n}\left(t, u_{1}, \ldots, u_{n}\right)
$$

- How to integrate in large dimensions ?
- Using a minimal number of evaluations of $f_{n}$ which costs $O\left(2^{n}\right)$


## Integration in large dimensions

Dimension
of the integral

$\mathbf{N}=$ Number of computed points of the integrand

Integration in large dimensions
J. Dick, F.Y. Kuo, I.H. Sloan
"High-dimensional integration:
The Quasi-Monte Carlo way," Acta Numerica 22, I 33 (2013).

Dimension
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$\mathbf{N}=$ Number of computed points of the integrand

## Quasi-Monte Carlo

- Evaluate the function on some special points. Low discrepancy sequences, e.g. Sobol'.
- Theorems

If the function $f$ is "smooth enough" (i.e. proper functional space), then


Sobol'sequence

Sobol'


Quasi-Monte Carlo IS NOT Monte Carlo. No random numbers.

Is our function smooth enough ?
No, but ...

- Change of variable $u(x)$ in $n$ dimensions.
- Goal : make the function flat/smooth

$$
Q_{n}=\int d^{n} \boldsymbol{u} f_{n}\left(u_{1}, \ldots, u_{n}\right) \quad \longrightarrow \quad Q_{n}=\int_{[0,1]^{n}} \mathrm{~d}^{n} \boldsymbol{x} f_{n}[\boldsymbol{u}(\boldsymbol{x})]\left|\frac{\partial \boldsymbol{u}}{\partial \boldsymbol{x}}\right|
$$

- $u(x)$ constructed from a model function $p_{n}(\mathbf{u})$ such that

$$
\left|\frac{\partial \boldsymbol{u}}{\partial \boldsymbol{x}}\right|=\frac{\mathcal{C}}{p_{n}(\boldsymbol{u})}
$$

- Then use quasi-Monte Carlo in new variable $x$.

$$
Q_{n} \approx Q_{n}(N)=\frac{\mathcal{C}}{N} \sum_{i=0}^{N} \frac{f_{n}\left[\boldsymbol{u}\left(\bar{x}_{i}\right)\right]}{p_{n}\left[\boldsymbol{u}\left(\bar{x}_{i}\right)\right]} \text { Sobol'sequence }
$$

## Model function

- Model function : approximation of the integrand in $n$ dimensions.

Optimized for quicker convergence.

- Machine learning problem.
- In general, a large class of possible functions, e.g. functional tensor trains / MPS

$$
p_{n}(\boldsymbol{u})=h_{a}^{(1)}\left(t-u_{1}\right) h_{a b}^{(2)}\left(u_{1}-u_{2}\right) \cdots h_{c d}^{(n-1)}\left(u_{n-2}-u_{n-1}\right) h_{d}^{(n)}\left(u_{n-1}-u_{n}\right)
$$

- Here, even the simplest case, without any optimization, already gives excellent results.

$$
p_{n}(\boldsymbol{u})=\prod_{i=1}^{n} h^{(i)}\left(u_{i-1}-u_{i}\right) \quad h^{(i)}(u)=e^{-u / \tau}
$$

## Quantum Quasi-Monte Carlo (QQMC)

- Compute the integral of the quantum problem with Quasi-Monte Carlo



## Error scaling with N

- Same curve in log-log



## Warping the integral is crucial



- Model function + quasi-MC = best method


## Large orders

- Error vs analytical Bethe Ansatz result, vs the number of sampling points $\mathbf{N}$



## Kondo ridge

## Experiment


T. Delattre et al. Nat. Phvs. 208 (2009)

Coulomb diamond $T=0$. vs $U$ and $\varepsilon_{d}$

M. Macek, P. Dumitrescu, C. Bertrand, B.Triggs, OP, X. Waintal ArXiv:2002.I 2372

- Many calculations ("parametric runs"), for various $U$ and $\varepsilon_{d}$.

About 25 cpu hours/point for order 10.

## Conclusion

- Solution of the out of equilibrium quantum dot.
- Perturbation theory even at strong coupling (with resummation)
- Quantum Quasi-Monte Carlo
- Roadmap : DMFT solvers, lattice problems ...

Thank you for your attention!

## Sobol' points

- Illustration in $d=2$

$$
N=100 \quad N=500 \quad N=1000
$$



## Learning the model function

- Optimization of model function from the data
- Proof of concepts : gain speed factor $\times 2$, using a projection technique


