# Hierarchical low rank compression, with an application in quantum many-body physics

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CCM/CCQ

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On work with Denis Golež of CCQ





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$$\begin{pmatrix} 2 & -1 & 0 & \cdots & 0 \\ -1 & 2 & \ddots & & \vdots \\ 0 & \ddots & \ddots & -1 & 0 \\ \vdots & & -1 & 2 & -1 \\ 0 & \cdots & 0 & -1 & 2 \end{pmatrix}$$

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  - FFT uses *recursive* structure of discrete Fourier transform matrix to get fast matrix-vector product:  $O(N^2) \rightarrow O(N \log N)$ 
    - Low rank and hierarchical low rank structure

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- Yes TSVD gives O(N) algorithm

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Targets = sources, with regularized interaction:

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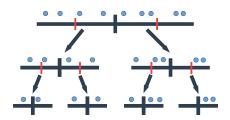
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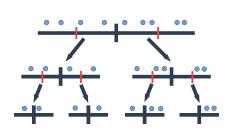
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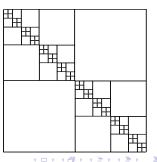
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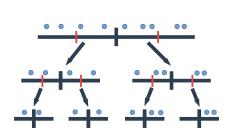


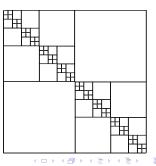
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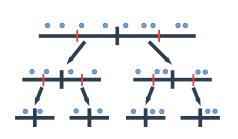


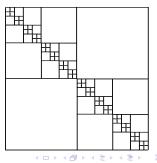




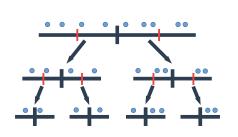


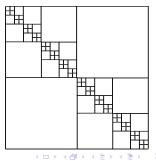
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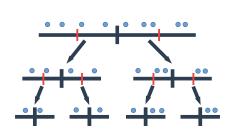


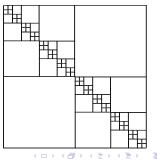
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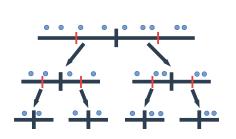
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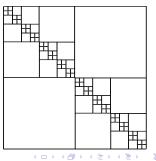




## Hierarchical off-diagonal low rank (HODLR) matrices

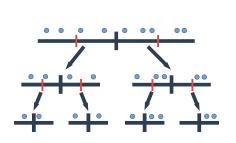
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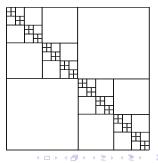




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- Want *k* fixed (or *slowly growing*) as  $N \to \infty$





# Application: nonequilibrium Dyson equation from quantum many-body physics

Joint work with Denis Golež of CCQ

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- Exponential complexity goes into interaction kernel (self-energy  $\Sigma$ ) between  $G_{jk}$ , which is truncated at some order

#### Nonequilibrium Dyson equations

$$(-\partial_{\tau} - h(0)) G^{M}(\tau) - \int_{0}^{\beta} d\bar{\tau} \, \Sigma^{M}(\tau - \bar{\tau}) G^{M}(\bar{\tau}) = 0 \tag{1}$$

$$(-i\partial_{t'} - h(t')) G^{R}(t, t') - \int_{t'}^{t} d\bar{t} G^{R}(t, \bar{t}) \Sigma^{R}(\bar{t}, t') = 0$$
 (2)

$$(i\partial_{t} - h(t)) G^{\uparrow}(t, \tau) - \int_{0}^{t} d\overline{t} \Sigma^{R}(t, \overline{t}) G^{\uparrow}(\overline{t}, \tau)$$

$$= \int_{0}^{\beta} d\overline{\tau} \Sigma^{\uparrow}(t, \overline{\tau}) G^{M}(\overline{\tau} - \tau)$$
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$$\left(\mathrm{i}\partial_t - h(t)\right)G^<(t,t') - \int_0^t d\bar{t} \, \Sigma^\mathrm{R}(t,\bar{t})G^<(\bar{t},t')$$

$$= \int_0^{t'} d\bar{t} \, \Sigma^{<}(t,\bar{t}) G^A(\bar{t},t') - \mathrm{i} \int_0^\beta d\bar{\tau} \, \Sigma^{\rceil}(t,\bar{\tau}) G_1^{\lceil}(\bar{\tau},t')$$

$$G^{M}(-\tau) = \xi G^{M}(\beta - \tau) \tag{5}$$

$$G^{R}(t,t) = -i$$

$$G^{\uparrow}(0,\tau) = iG^{M}(-\tau) = i\xi G^{M}(\beta - \tau)$$

$$G^{<}(0,t') = -\overline{G^{\uparrow}(t',0)}$$

$$G^{\lceil}(\tau, t) = -\varepsilon \overline{G^{\rceil}(t, \beta - \tau)}$$

$$G^{A}(t,t') = \overline{G^{R}(t',t)}$$

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(4)

(6)

(7)

(8)

(9)

(10)

(11)

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Focus on this 
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$$\frac{f(t,t) - \zeta_{G}(t,p-t)}{t}$$

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For each fixed t', after some manipulations...

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$$y'(t) + p(t)y(t) + \int_0^t k(t, s)y(s) ds = 0$$
  
 $y(0) = y_0$   
 $t \in [0, T]$ 

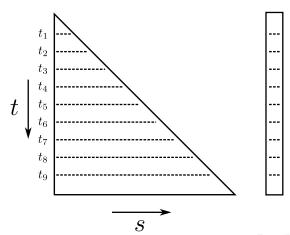
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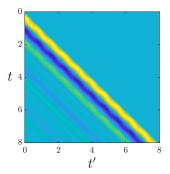
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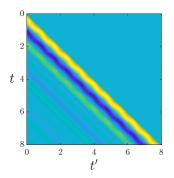


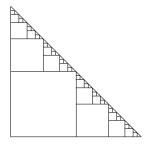
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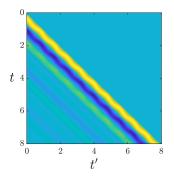


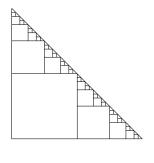
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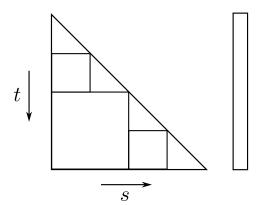




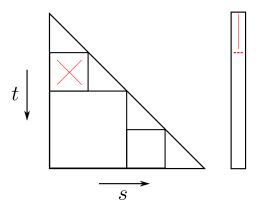
ullet Have efficient method of building HODLR representation of  $\Sigma^R(t,t')$ 

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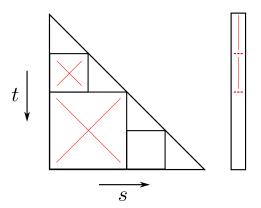
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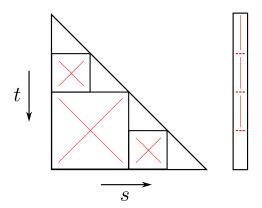
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